Time Domain Maxwell's Equations

starting with Maxwell's curl equations

$$\bar{\nabla} \times \bar{E} = -\frac{\partial \bar{B}}{\partial t} \qquad \bar{\nabla} \times \bar{H} = \frac{\partial \bar{D}}{\partial t} + \bar{J}$$

in a source-free, homogeneous, isotropic, linear, and stationary medium:

$$\bar{\nabla} \times \bar{E} = -\mu \frac{\partial \bar{H}}{\partial t} \qquad \bar{\nabla} \times \bar{H} = \epsilon \frac{\partial \bar{E}}{\partial t}$$

which represent 6 scalar equations in 6 scalar unknowns.

In rectangular coordinates these can be written out component-wize as

$$\bar{\nabla} \times \bar{A} = (\partial_y A_z - \partial_z A_y)\hat{a}_x + (\partial_z A_x - \partial_x A_z)\hat{a}_y + (\partial_x A_y - \partial_y A_x)\hat{a}_z$$

$$\begin{cases} -\mu\, \partial_t H_x = \partial_y E_z - \partial_z E_y \\ -\mu\, \partial_t H_y = \partial_z E_x - \partial_x E_z \\ -\mu\, \partial_t H_z = \partial_x E_y - \partial_y E_x \end{cases}$$

$$\begin{cases} \epsilon\, \partial_t E_x = \partial_y H_z - \partial_z H_y \\ \epsilon\, \partial_t E_y = \partial_z H_x - \partial_x H_z \\ \epsilon\, \partial_t E_z = \partial_x H_y - \partial_y H_x \end{cases}$$

these can be written in general vector form as

$$A\,\partial_t \underline{u} + \partial_x \underline{E} + \partial_y \underline{F} + \partial_z \underline{G} = \underline{0}$$

$$\underline{u} = \begin{pmatrix} E_x \\ E_y \\ E_z \\ H_x \\ H_y \\ H_z \end{pmatrix} \qquad \underline{E} = \begin{pmatrix} 0 \\ H_z \\ -H_y \\ 0 \\ -E_z \\ E_y \end{pmatrix} \qquad \underline{F} = \begin{pmatrix} -H_z \\ 0 \\ H_x \\ E_z \\ 0 \\ -E_x \end{pmatrix} \qquad \underline{G} = \begin{pmatrix} H_y \\ -H_x \\ 0 \\ -E_y \\ E_x \\ 0 \end{pmatrix}$$

$$A = \mathrm{diag}\{\epsilon\ \epsilon\ \epsilon\ \mu\ \mu\ \mu\} = \begin{bmatrix} \epsilon & 0 & 0 & 0 & 0 & 0 \\ 0 & \epsilon & 0 & 0 & 0 & 0 \\ 0 & 0 & \epsilon & 0 & 0 & 0 \\ 0 & 0 & 0 & \mu & 0 & 0 \\ 0 & 0 & 0 & 0 & \mu & 0 \\ 0 & 0 & 0 & 0 & 0 & \mu \end{bmatrix}$$

$$A^{-1} = \mathrm{diag}\{e\ e\ e\ m\ m\ m\} \qquad e = \frac{1}{\epsilon} \qquad m = \frac{1}{\mu}$$

$$\therefore \quad \boxed{\partial_t \underline{u} + A^{-1}\partial_x \underline{E} + A^{-1}\partial_y \underline{F} + A^{-1}\partial_z \underline{G} = \underline{0}}$$

One dimensional waves:

if we have no variation of $\underline{u}$ with respect to the y and z directions then $\partial_y = \partial_z = 0$ and our system of equations becomes:

$$\partial_t \underline{u} + A^{-1} \partial_x \underline{E} = \underline{0}$$

$$\partial_t \begin{pmatrix} E_x \\ E_y \\ E_z \\ H_x \\ H_y \\ H_z \end{pmatrix} + A^{-1} \partial_x \begin{pmatrix} 0 \\ H_z \\ -H_y \\ 0 \\ -E_z \\ E_y \end{pmatrix} = \underline{0}$$

From this we see that we have 4 independent systems:

$$\partial_t E_x = 0 \qquad (1)$$

$$\partial_t H_x = 0 \qquad (2)$$

$$\left.\begin{aligned} \partial_t E_y + e\, \partial_x H_z &= 0 \\ \partial_t H_z + m\, \partial_x E_y &= 0 \end{aligned}\right\} \qquad (3)$$

$$\left.\begin{aligned} \partial_t E_z - e\, \partial_x H_y &= 0 \\ \partial_t H_y - m\, \partial_x E_z &= 0 \end{aligned}\right\} \qquad (4)$$

Solutions of ① + ② :

$$\partial_t \psi(x,t) = 0 \quad \Rightarrow \quad \psi = \psi(x)$$

$$\therefore \begin{cases} E_x(x,t) = E_0(x) & \text{(initial condition)} \\ H_x(x,t) = H_0(x) & \text{(initial condition)} \end{cases}$$

"static" fields , no variations in time

<u>solution of ③</u> :

$$\boxed{\partial_t \begin{pmatrix} E_y \\ H_z \end{pmatrix} + \begin{bmatrix} 0 & e \\ m & 0 \end{bmatrix} \partial_x \begin{pmatrix} E_y \\ H_z \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \end{pmatrix}}$$

Finite Difference Methods for Time Domain Maxwell's equations

starting with the 1-D system:

$$\left.\begin{array}{l} \partial_t \underline{u} + A \partial_x \underline{u} = \underline{0} \\[1em] \underline{u} = \begin{pmatrix} E_y \\ H_z \end{pmatrix} \qquad A = \begin{bmatrix} 0 & e \\ m & 0 \end{bmatrix} \end{array}\right\}$$

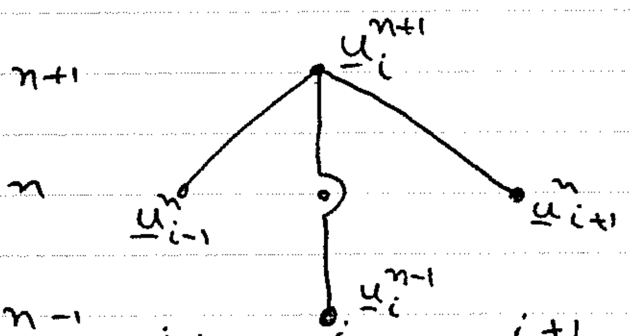
use the center-difference approximation to the partial derivative operators

$$\partial_t \underline{u} = \frac{\underline{u}_i^{n+1} - \underline{u}_i^{n-1}}{2\Delta t} + O(\Delta t^2)$$

$$\partial_x \underline{u} = \frac{\underline{u}_{i+1}^{n} - \underline{u}_{i-1}^{n}}{2\Delta x} + O(\Delta x^2)$$

$$\frac{\underline{u}_i^{n+1} - \underline{u}_i^{n-1}}{2\Delta t} + A \frac{\underline{u}_{i+1}^{n} - \underline{u}_{i-1}^{n}}{2\Delta x} = \underline{0}$$

$$\boxed{\underline{u}_i^{n+1} = \underline{u}_i^{n-1} - \frac{\Delta t}{\Delta x} A \left(\underline{u}_{i+1}^{n} - \underline{u}_{i-1}^{n} \right)}$$



" computational " molecule

the above is called <u>Leap - Frog</u> scheme

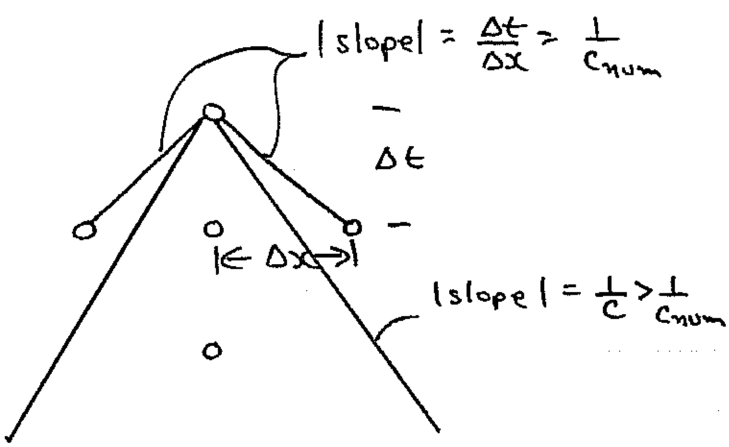
and is a two - step scheme since we require $\underline{u}^0$ and $\underline{u}^1$ to start the scheme.

It can be shown that this scheme is stable for:

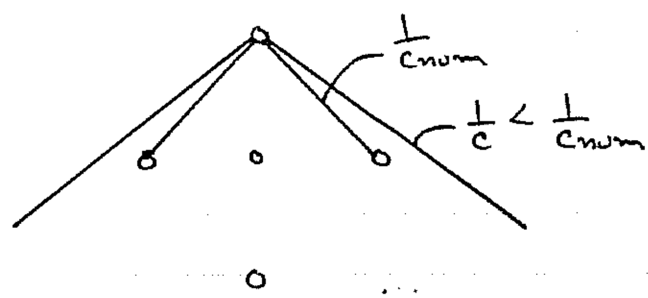
$$\boxed{c\frac{\Delta t}{\Delta x} \leq 1}$$ Courant - Friedrichs - Lewy Condition (CFL)

$$\frac{\Delta x}{\Delta t} = c_{num} \quad \text{— numerical speed}$$

$\therefore$ CFL condition implies that the numerical speed must be greater than the actual speed.



Stable
$\underline{c_{num} > c}$



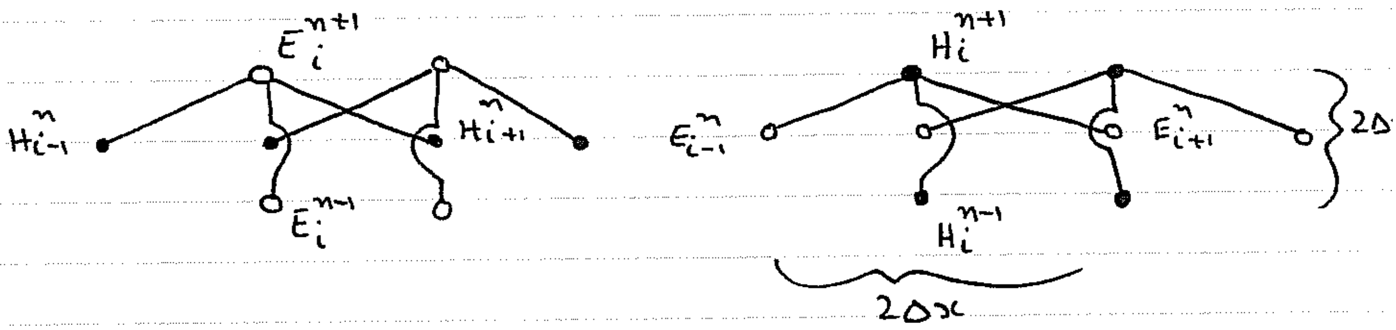
Unstable.
$\underline{c > c_{num}}$

now writting these component wize

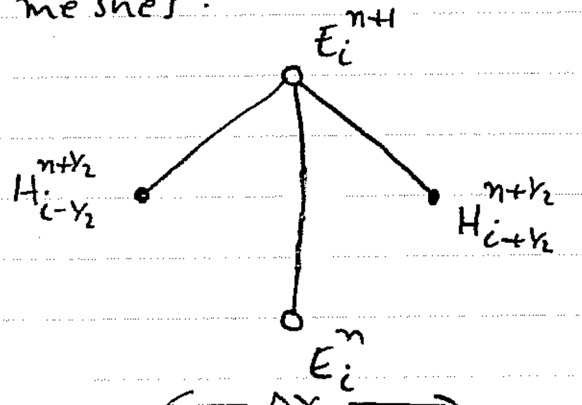
let $\begin{pmatrix} E_y \\ H_z \end{pmatrix} = \begin{pmatrix} E \\ H \end{pmatrix}$

$$\left.\begin{aligned} E_i^{n+1} &= E_i^{n-1} - \frac{\Delta t}{\Delta x}\, e\left(H_{i+1}^{n} - H_{i-1}^{n}\right) \\ H_i^{n+1} &= H_i^{n-1} - \frac{\Delta t}{\Delta x}\, m\left(E_{i+1}^{n} - E_{i-1}^{n}\right) \end{aligned}\right\}$$

<u>computational molecules</u>:



a close look reveals <u>4 independent</u> meshes which are interlaced in space and time.
we only need to keep one of these meshes:



Δt

$$\frac{\Delta x}{\Delta t} > c$$

$$\boxed{\begin{aligned} E_i^{n+1} &= E_i^{n} - \frac{\Delta t}{\Delta x} e \left(H_{i+1/2}^{n+1/2} - H_{i-1/2}^{n+1/2} \right) \\ H_{i+1/2}^{n+1/2} &= H_{i+1/2}^{n-1/2} - \frac{\Delta t}{\Delta x} m \left(E_{i+1}^{n} - E_i^{n} \right) \end{aligned}}$$

1-D Yee Algorithm.

Given an initial $\bar{E}$ field $-E(x,0) = E_0(x)$

we set $E_i^0 = E_0(i\Delta x)$

the first set of H fields are then computed as.

$$\boxed{H_{i+1/2}^{n+1/2} = \frac{-\Delta t}{2\Delta x} m \left(E_{i+1}^{n} - E_i^{n} \right)}$$

if we choose $\frac{\Delta t}{\Delta x} = \frac{1}{c} = (em)^{-1/2}$

then $\frac{\Delta t}{\Delta x} e = e(em)^{-1/2} = \sqrt{\frac{e}{m}} = \sqrt{\frac{\mu}{\epsilon}} = Z$

$\frac{\Delta t}{\Delta x} m = m(em)^{-1/2} = \sqrt{\frac{m}{e}} = \sqrt{\frac{\epsilon}{\mu}} = Y = \frac{1}{Z}$

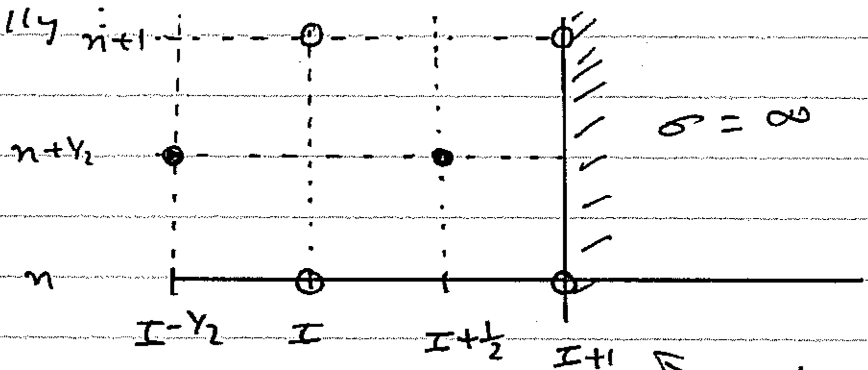
in free space $Z \approx 376.734 \approx 376.7$

$(\mu_0 = 4\pi \times 10^{-7} \quad \epsilon_0 = 8.854 \times 10^{-12})$

Implementing Boundary Conditions & Inhomogeneous Medium

1-D Yee:

at a perfectly conducting we know that the tangential electric field must be zero. Since the magnetic field is specified at interlaced points in space they can be handled normally

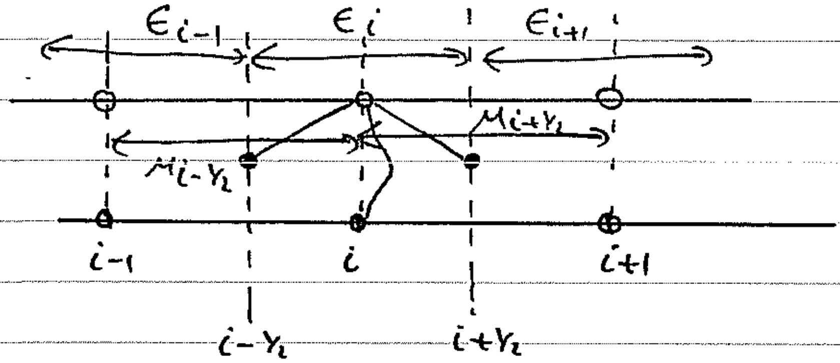


end of mesh. Since nothing gets through

$$\begin{cases} E_{I+1}^{n} = 0 \\ H_{I+1/2}^{n+1/2} = H_{I+1/2}^{n-1/2} - \dfrac{\Delta t}{\Delta x}\, m_{I+1/2} \left(\underbrace{E_{I+1}^{n}}_{0} - E_{I}^{n} \right) \end{cases}$$

∴ Rule is simple:

Put your perfect conducting boundary at an "E" point and keep E = 0 at that point inhomogeneous μ and ϵ are also simply handled :



$$\begin{cases} E_i^{n+1} = E_i^n - \dfrac{\Delta t}{\Delta x}\, e_i \left(H_{i+1/2}^{n+1/2} - H_{i-1/2}^{n+1/2} \right) \\[2ex] H_{i+1/2}^{n+1/2} = H_{i+1/2}^{n-1/2} - \dfrac{\Delta t}{\Delta x}\, m_{i+1/2} \left(E_{i+1}^n - E_i^n \right) \end{cases}$$

unfortunately boundaries of μ and boundaries of ϵ <u>cannot</u> coincide !

Non-pefectly conducting Media

$\sigma \neq 0$

$$\bar{\nabla} \times \bar{E} = -\frac{\partial \bar{B}}{\partial t} = -\mu \frac{\partial \bar{H}}{\partial t}$$

Ohms Law ($\bar{J} = \sigma \bar{E}$)

$$\bar{\nabla} \times \bar{H} = \frac{\partial \bar{D}}{\partial t} + \bar{J} = \epsilon \frac{\partial E}{\partial t} + \sigma \bar{E}$$

in 1-D the curl equations become:

$$\begin{cases} \partial_t \underline{u} + A \partial_x \underline{u} = \underline{s} \\ \underline{u} = \begin{pmatrix} E_y \\ H_z \end{pmatrix} = \begin{pmatrix} E \\ H \end{pmatrix} \qquad A = \begin{pmatrix} 0 & e \\ m & 0 \end{pmatrix} \qquad \underline{s} = \begin{pmatrix} -\sigma e E \\ 0 \end{pmatrix} \end{cases}$$

Show this!

Now how do we difference this equation?

$$\frac{\underline{u}_i^{n+1} - \underline{u}_i^{n-1}}{2\Delta t} + A \frac{\underline{u}_{i+1}^n - \underline{u}_{i-1}^n}{2\Delta x} = \underline{s}_i^{n+1}$$

must use $n+1$ for stability.

using Yee version component wize:

$$E_i^{n+1} + \Delta t \sigma e E_i^{n+1} = E_i^n - \frac{\Delta t}{\Delta x} e_i \left(H_{i+1/2}^{n+1/2} - H_{i-1/2}^{n+1/2} \right)$$

$$E_i^{n+1} = \left(\frac{1}{1 + \Delta t \sigma_i e_i} \right) \left[E_i^n - \frac{\Delta t}{\Delta x} e_i \left(H_{i+1/2}^{n+1/2} - H_{i-1/2}^{n+1/2} \right) \right]$$

$H_{i+1/2}^{n+1/2}$ — same as before.

2-D FIELDS. $\partial_z = 0$

$$\partial_t \underline{u} + A^{-1}\partial_x \underline{E} + A^{-1}\partial_y \underline{F} = \underline{0}$$

$$\partial_t \begin{pmatrix} E_x \\ E_y \\ E_z \\ H_x \\ H_y \\ H_z \end{pmatrix} + A^{-1}\partial_x \begin{pmatrix} 0 \\ H_z \\ -H_y \\ 0 \\ -E_z \\ E_y \end{pmatrix} + A^{-1}\partial_y \begin{pmatrix} -H_z \\ 0 \\ H_x \\ E_z \\ 0 \\ -E_x \end{pmatrix} = \underline{0}$$

we notice that there are two independent systems of equation

Transverse Electric to z (TE_z) (E_x, E_y, H_z)

$$\partial_t \begin{pmatrix} E_x \\ E_y \\ H_z \end{pmatrix} + \begin{bmatrix} e & 0 & 0 \\ 0 & e & 0 \\ 0 & 0 & m \end{bmatrix} \partial_x \begin{pmatrix} 0 \\ H_z \\ E_y \end{pmatrix} + \begin{bmatrix} e & 0 & 0 \\ 0 & e & 0 \\ 0 & 0 & m \end{bmatrix} \partial_y \begin{pmatrix} -H_z \\ 0 \\ -E_x \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \\ 0 \end{pmatrix}$$

Transverse Magnetic to z (TM_z) (H_x, H_y, E_z)

$$\partial_t \begin{pmatrix} E_z \\ H_x \\ H_y \end{pmatrix} + \begin{bmatrix} e & 0 & 0 \\ 0 & m & 0 \\ 0 & 0 & m \end{bmatrix} \partial_x \begin{pmatrix} -H_y \\ 0 \\ -E_z \end{pmatrix} + \begin{bmatrix} e & 0 & 0 \\ 0 & m & 0 \\ 0 & 0 & m \end{bmatrix} \partial_y \begin{pmatrix} H_x \\ E_z \\ 0 \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \\ 0 \end{pmatrix}$$

we can now apply the Leap-Frog scheme as :

$$\frac{\underline{u}_{ij}^{n+1} - \underline{u}_{ij}^{n-1}}{2\Delta t} + \vec{A}\left(\frac{\underline{E}_{i+1j}^{n} - \underline{E}_{i-1j}^{n}}{2\Delta x} + \frac{\underline{F}_{ij+1}^{n} - \underline{F}_{ij-1}^{n}}{2\Delta y}\right) = \underline{0}$$

$$\underline{u}_{ij}^{n+1} = \underline{u}_{ij}^{n-1} - \frac{\Delta t A'}{\Delta x}\left(\underline{E}_{i+1j}^{n} - \underline{E}_{i-1j}^{n}\right) - \frac{\Delta t}{\Delta y} A'\left(\underline{F}_{ij+1}^{n} - \underline{F}_{ij-1}^{n}\right)$$

let $\Delta x = \Delta y = h$, $\Delta t = k$

$$\boxed{\underline{u}_{ij}^{n+1} = \underline{u}_{ij}^{n-1} - \frac{k}{h} A'\left(\underline{E}_{i+1j}^{n} - \underline{E}_{i-1j}^{n} + \underline{F}_{ij+1}^{n} - \underline{F}_{ij-1}^{n}\right)}$$

specific A' , $\underline{E}$, $\underline{F}$ depends on TE_z , or TM_z.

This scheme is stable for (3-D)

$$\Delta t \le \frac{1}{\left(\frac{1}{\Delta x^2} + \frac{1}{\Delta y^2} + \frac{1}{\Delta z^2}\right)^{1/2} c_{max}}$$

in 2-D , $\Delta x = \Delta y = \Delta h$

$$\boxed{\frac{\Delta t}{\Delta h} c_{max} \le \frac{1}{\sqrt{2}}}$$

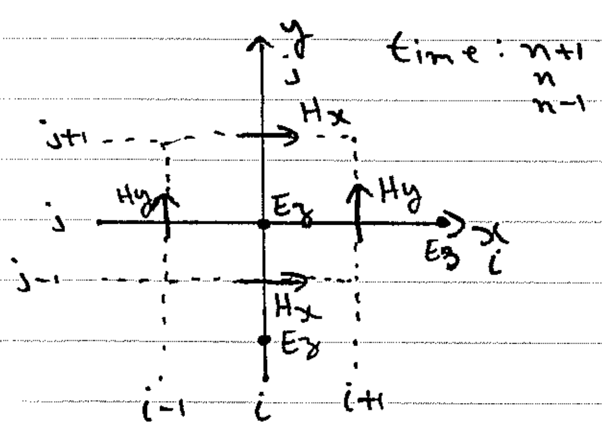
Application of the Leap-Frog Formula:

TM waves: $(\Delta x = \Delta y = h \qquad \Delta t = k)$

$$\underline{u} = \begin{pmatrix} E_z \\ H_x \\ H_y \end{pmatrix} \qquad \underline{E} = \begin{pmatrix} -H_y \\ 0 \\ -E_z \end{pmatrix} \qquad \underline{F} = \begin{pmatrix} H_x \\ E_z \\ 0 \end{pmatrix}$$

$$A^{-1} = \begin{bmatrix} e & 0 & 0 \\ 0 & m & 0 \\ 0 & 0 & m \end{bmatrix}$$

$$\begin{cases} E_{z_{ij}}^{n+1} = E_{z_{ij}}^{n-1} - \dfrac{k e_{ij}}{h}\left(-H_{y_{i+1j}}^{n} + H_{y_{i-1j}}^{n} + H_{x_{ij+1}}^{n} - H_{x_{ij-1}}^{n}\right) \\[2ex] H_{x_{ij}}^{n+1} = H_{x_{ij}}^{n-1} - \dfrac{k}{h} m_{ij}\left(E_{z_{ij+1}}^{n} - E_{z_{ij-1}}^{n}\right) \\[2ex] H_{y_{ij}}^{n+1} = H_{y_{ij}}^{n-1} - \dfrac{k}{h} m_{ij}\left(-E_{z_{i+1j}}^{n} + E_{z_{i-1j}}^{n}\right) \end{cases}$$



"computational molecule."

It is a bit harder to see now but this scheme is also interlaced in both time and space.

example: E_z at even time points $n = 0, 2, 4, \ldots$

depend on:

other E_z at even time points and H_x, H_y at odd time points.

this results in independent interlaced meshes.

Keeping only one mesh:

E_z at $n = 0, 1, 2, \ldots$

$i, j = \pm 1, \pm 2, \ldots$

H_x at $n = \frac{1}{2}, 1+\frac{1}{2}, 2+\frac{1}{2}$ $\quad n+\frac{1}{2}$

$i = \pm 1, \pm 2, \ldots$ $\quad i$

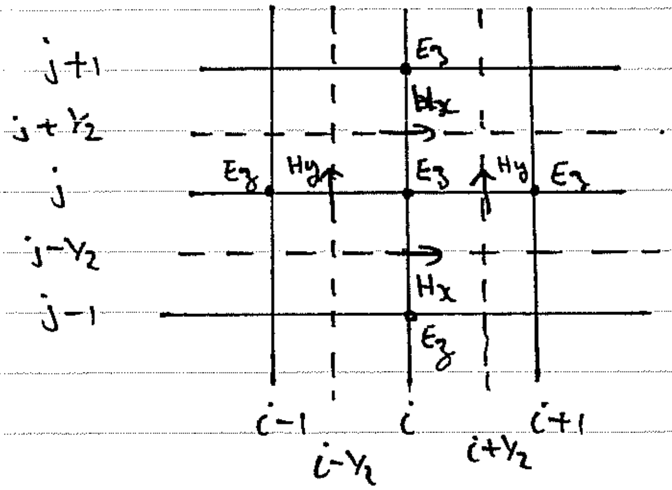
$j = \pm \frac{1}{2}, \pm(1+\frac{1}{2}), \pm(2\pm\frac{1}{2})$ $\quad j+\frac{1}{2}$

H_y at $n = \frac{1}{2}, 1+\frac{1}{2}, \ldots$ $\quad n+\frac{1}{2}.$

$i = \pm\frac{1}{2}, \ldots$ $\quad i+\frac{1}{2}$

$j = \pm 1, \pm 2$ $\quad j$

now the computational molecule becomes.



(Yee Algorithm)

$$
\begin{cases}
E_{z_{ij}}^{n+1} = E_{z_{ij}}^{n} - \frac{k}{h} e_{ij} \left(-H_{y_{i+\frac{1}{2} j}}^{n+\frac{1}{2}} + H_{y_{i-\frac{1}{2} j}}^{n+\frac{1}{2}} + H_{x_{i j+\frac{1}{2}}}^{n+\frac{1}{2}} - H_{x_{i j-\frac{1}{2}}}^{n+\frac{1}{2}} \right) \\
H_{x_{i j+\frac{1}{2}}}^{n+\frac{1}{2}} = H_{x_{i j+\frac{1}{2}}}^{n-\frac{1}{2}} - \frac{k}{h} m_{i j+\frac{1}{2}} \left(E_{z_{i j+1}}^{n} - E_{z_{ij}}^{n} \right) \\
H_{y_{i+\frac{1}{2} j}}^{n+\frac{1}{2}} = H_{y_{i+\frac{1}{2} j}}^{n-\frac{1}{2}} - \frac{k}{h} m_{i+\frac{1}{2} j} \left(-E_{z_{i+1 j}}^{n} + E_{z_{ij}}^{n} \right)
\end{cases}
$$

We would initialize the E_{ij}^{0} with the initial conditions and $H^{-\frac{1}{2}} = 0$ or I.C. if they are known.

Full 3-D plus Time Difference Methods

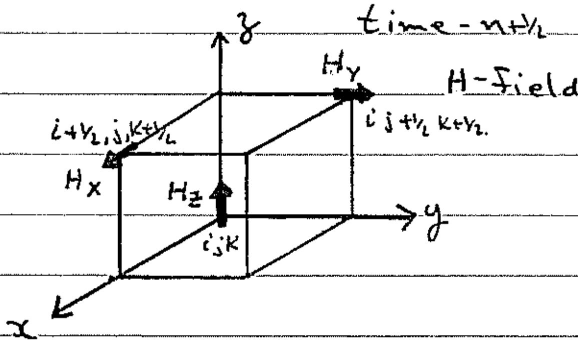
starting with the conservation law:

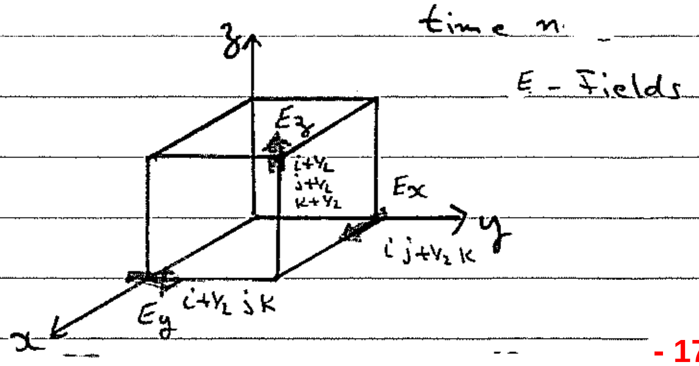
$$\partial_t \underline{u} + A^{-1}\partial_x \underline{E} + A^{-1}\partial_y \underline{F} + A^{-1}\partial_z \underline{G} = \underline{0}$$

expanding:

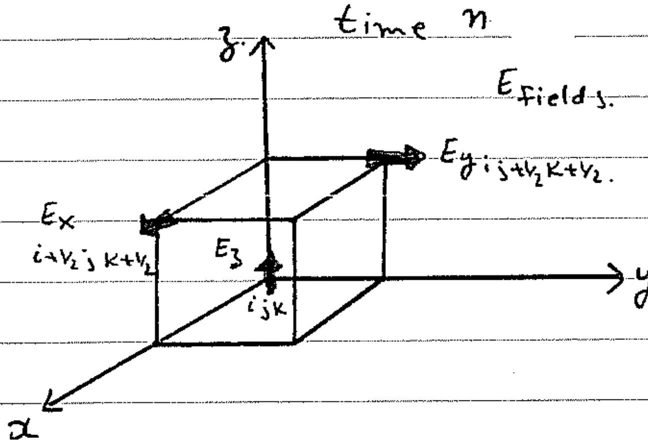
$$\begin{cases} \partial_t E_x + e\,\partial_x(0) + e\,\partial_y(-H_z) + e\,\partial_z(H_y) = 0 \\ \partial_t E_y + e\,\partial_x(H_z) + e\,\partial_y(0) + e\,\partial_z(-H_x) = 0 \\ \partial_t E_z + e\,\partial_x(-H_y) + e\,\partial_y(H_x) + e\,\partial_z(0) = 0 \\ \partial_t H_x + m\,\partial_x(0) + m\,\partial_y(E_z) + m\,\partial_z(-E_y) = 0 \\ \partial_t H_y + m\,\partial_x(-E_z) + m\,\partial_y(0) + m\,\partial_z(E_x) = 0 \\ \partial_t H_z + m\,\partial_x(E_y) + m\,\partial_y(-E_x) + m\,\partial_z(0) = 0 \end{cases}$$

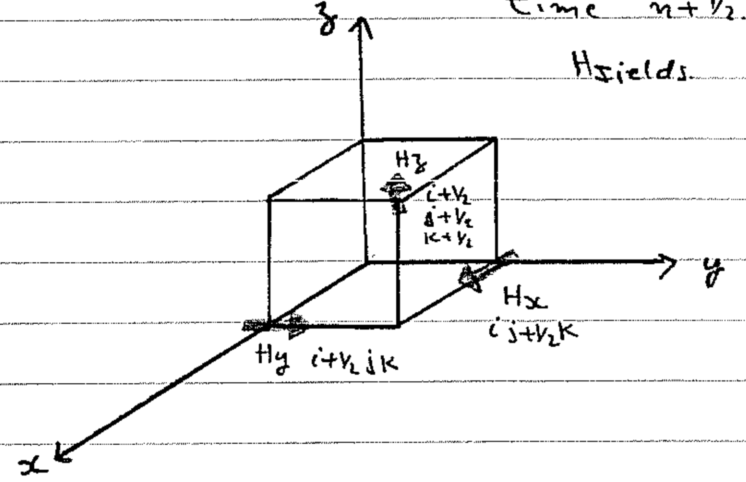
if we difference these equations using the leap-frog method (i.e. centered time and centered space) then again we end up with independent meshes which are interlaced in space and time. We retain only one mesh by "putting" E's on integer time steps and H's on $1/2$ time steps





alternatively :





Note: in first set:

bottom of cubes : TE_z.

top of cubes : TM_z.

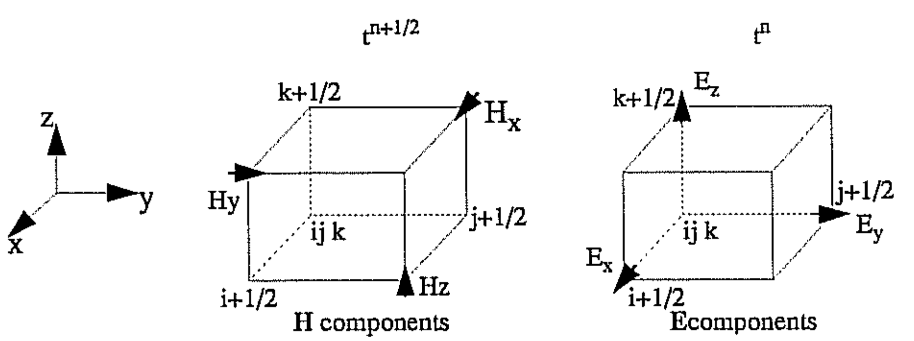
in second set :

bottom of cubes : TM_z.

top of cubes : TE_z.

it doesn't really make any difference whether we choose one set or the other

a set which will simplify notation is a slight modification of the first:



say we take this last set and formulate the leap-frog scheme with
$\Delta x = \Delta y = \Delta z = h \qquad \Delta t = k$:

$$
\left\{
\begin{aligned}
E_{x\,i+\frac12\, j\, k}^{\,n+1} &= E_{x\,i+\frac12\, j\, k}^{\,n} - \frac{k\,e}{h}\Big(H_{z\,i+\frac12\, j-\frac12\, k}^{\,n+\frac12} - H_{z\,i+\frac12\, j+\frac12\, k}^{\,n+\frac12} + \\
&\qquad\qquad H_{y\,i+\frac12\, j\, k+\frac12}^{\,n+\frac12} - H_{y\,i+\frac12\, j\, k-\frac12}^{\,n+\frac12} \Big) \\
E_{y\,i\, j+\frac12\, k}^{\,n+1} &= E_{y\,i\, j+\frac12\, k}^{\,n} - \frac{k\,e}{h}\Big(H_{z\,i+\frac12\, j+\frac12\, k}^{\,n+\frac12} - H_{z\,i-\frac12\, j+\frac12\, k}^{\,n+\frac12} + \\
&\qquad\qquad H_{x\,i\, j+\frac12\, k-\frac12}^{\,n+\frac12} - H_{x\,i\, j+\frac12\, k+\frac12}^{\,n+\frac12} \Big) \\
E_{z\,i\, j\, k+\frac12}^{\,n+1} &= E_{z\,i\, j\, k+\frac12}^{\,n} - \frac{k\,e}{h}\Big(H_{y\,i-\frac12\, j\, k+\frac12}^{\,n+\frac12} - H_{y\,i+\frac12\, j\, k+\frac12}^{\,n+\frac12} + \\
&\qquad\qquad H_{x\,i\, j+\frac12\, k+\frac12}^{\,n+\frac12} - H_{x\,i\, j-\frac12\, k+\frac12}^{\,n+\frac12} \Big)
\end{aligned}
\right.
$$

$$
\left\{
\begin{aligned}
H_{x\,i\, j+\frac12\, k+\frac12}^{\,n+\frac12} &= H_{x\,i\, j+\frac12\, k+\frac12}^{\,n-\frac12} - \frac{k\,m}{h}\Big(E_{z\,i\, j+1\, k+\frac12}^{\,n} - E_{z\,i\, j\, k+\frac12}^{\,n} + \\
&\qquad\qquad E_{y\,i\, j+\frac12\, k}^{\,n} - E_{y\,i\, j+\frac12\, k+1}^{\,n} \Big) \\
H_{y\,i+\frac12\, j\, k+\frac12}^{\,n+\frac12} &= H_{y\,i+\frac12\, j\, k+\frac12}^{\,n-\frac12} - \frac{k\,m}{h}\Big(E_{z\,i\, j\, k+\frac12}^{\,n} - E_{z\,i+1\, j\, k+\frac12}^{\,n} + \\
&\qquad\qquad E_{x\,i+\frac12\, j\, k+1}^{\,n} - E_{x\,i+\frac12\, j\, k}^{\,n} \Big) \\
H_{z\,i+\frac12\, j+\frac12\, k}^{\,n+\frac12} &= H_{z\,i+\frac12\, j+\frac12\, k}^{\,n-\frac12} - \frac{k\,m}{h}\Big(E_{y\,i+1\, j+\frac12\, k}^{\,n} - E_{y\,i\, j+\frac12\, k}^{\,n} + \\
&\qquad\qquad E_{x\,i+\frac12\, j\, k}^{\,n} - E_{x\,i+\frac12\, j+1\, k}^{\,n} \Big)
\end{aligned}
\right.
$$

since the $\vec{E}$ and $\vec{H}$ fields exist at different time steps (i.e. n and $n+\frac{1}{2}$) we can reduce the complexity of the above equations by introducing some new notation.

We first split the conservation law into two systems:

$$\begin{cases} \partial_t \underline{u} + A\,\partial_x \underline{E}(\underline{v}) + A\,\partial_y \underline{F}(\underline{v}) + A\,\partial_z \underline{G}(\underline{v}) = \underline{0} \\ \partial_t \underline{v} + B\,\partial_x \underline{E}(\underline{u}) + B\,\partial_y \underline{F}(\underline{u}) + B\,\partial_z \underline{G}(\underline{u}) = \underline{0} \end{cases}$$

$$\begin{cases} \underline{u} = \begin{pmatrix} E_x \\ E_y \\ E_z \end{pmatrix} \quad \underline{E}(\underline{v}) = \underline{E} = \begin{pmatrix} 0 \\ H_z \\ -H_y \end{pmatrix} \quad \underline{F}(\underline{v}) = \underline{F} = \begin{pmatrix} -H_z \\ 0 \\ H_x \end{pmatrix} \quad \underline{G}(\underline{v}) = \underline{G} = \begin{pmatrix} H_y \\ -H_x \\ 0 \end{pmatrix} \\ \underline{v} = \begin{pmatrix} H_x \\ H_y \\ H_z \end{pmatrix} \quad \underline{E}(\underline{u}) = \underline{H} = \begin{pmatrix} 0 \\ -E_z \\ E_y \end{pmatrix} \quad \underline{F}(\underline{u}) = \underline{I} = \begin{pmatrix} E_z \\ 0 \\ -E_x \end{pmatrix} \quad \underline{G}(\underline{u}) = \underline{J} = \begin{pmatrix} -E_y \\ E_x \\ 0 \end{pmatrix} \end{cases}$$

$$A = \operatorname{diag}\{e\ e\ e\} \qquad B = \operatorname{diag}\{m\ m\ m\}$$

discretized notation becomes:

$$\underline{u}^n_{ijk} = \begin{pmatrix} E_x{}^n_{i+\frac{1}{2}\,j\,k} \\ E_y{}^n_{i\,j+\frac{1}{2}\,k} \\ E_z{}^n_{i\,j\,k+\frac{1}{2}} \end{pmatrix} \qquad \underline{v}^n_{ijk} = \begin{pmatrix} H_x{}^{n+\frac{1}{2}}_{i\,j+\frac{1}{2}\,k+\frac{1}{2}} \\ H_y{}^{n+\frac{1}{2}}_{i+\frac{1}{2}\,j\,k+\frac{1}{2}} \\ H_z{}^{n+\frac{1}{2}}_{i+\frac{1}{2}\,j+\frac{1}{2}\,k} \end{pmatrix}$$

$$\underline{E}^n_{ijk} = \underline{E}(\underline{v}^n_{ijk}) \qquad \underline{F}^n_{ijk} = \underline{F}(\underline{v}^n_{ijk}) \qquad \underline{G}^n_{ijk} = \underline{G}(\underline{v}^n_{ijk})$$

$$\underline{H}^n_{ijk} = \underline{E}(\underline{u}^n_{ijk}) \qquad \underline{I}^n_{ijk} = \underline{F}(\underline{u}^n_{ijk}) \qquad \underline{J}^n_{ijk} = \underline{G}(\underline{u}^n_{ijk})$$

The update procedure is a <u>two-step form</u>:

$$\begin{cases} \underline{u}^{n+1}_{ijk} = \underline{u}^{n}_{ijk} - A\dfrac{k}{h}\left(\underline{E}^{n}_{ijk} - \underline{E}^{n}_{i-1jk} + \underline{F}^{n}_{ijk} - \underline{F}^{n}_{ij-1k} + \underline{G}^{n}_{ijk} - \underline{G}^{n}_{ijk-1}\right) \\ \underline{v}^{n+1}_{ijk} = \underline{v}^{n}_{ijk} - B\dfrac{k}{h}\left(\underline{H}^{n+1}_{i+1jk} - \underline{H}^{n+1}_{ijk} + \underline{I}^{n+1}_{ij+1k} - \underline{I}^{n+1}_{ijk} + \underline{J}^{n+1}_{ijk+1} - \underline{J}^{n+1}_{ijk}\right) \end{cases}$$

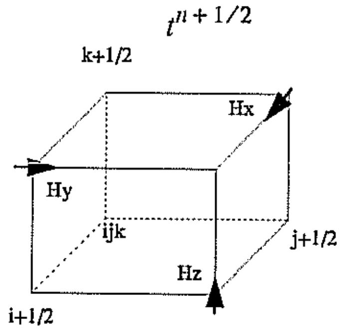
or

$$\begin{cases} \underline{u}^{n+1} = \underline{u}^{n} - \dfrac{k}{h} A\left(\nabla_x \underline{E}^{n} + \nabla_y \underline{F}^{n} + \nabla_z \underline{G}^{n}\right) \\ \underline{v}^{n+1} = \underline{v}^{n} - \dfrac{k}{h} B\left(\Delta_x \underline{H}^{n+1} + \Delta_y \underline{I}^{n+1} + \Delta_z \underline{J}^{n+1}\right) \end{cases}$$

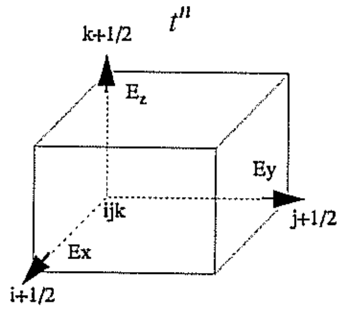
where $\nabla \rightarrow$ backward difference operator
$\Delta \rightarrow$ forward difference operator

eg: $\nabla_x \underline{E}^n = \underline{E}^n_{ijk} - \underline{E}^n_{i-1jk}$

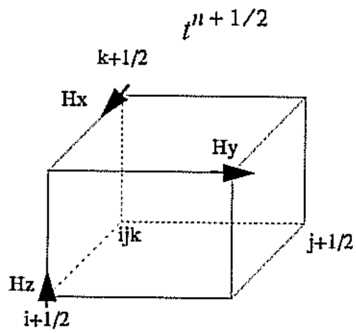
The Leap-frog scheme can be applied to the Maxwell curl equations and the result is sixteen independent sets of space-time interleaved discretized electric and magnetic fields. Eight of these are shown below. In order to get the remaining 8 the time interlacing between electric and magnetic fields is simply interchanged.



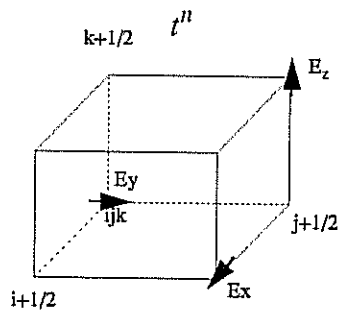
H components



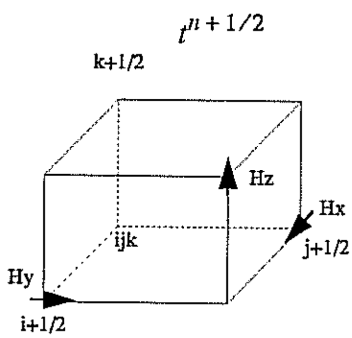
E components



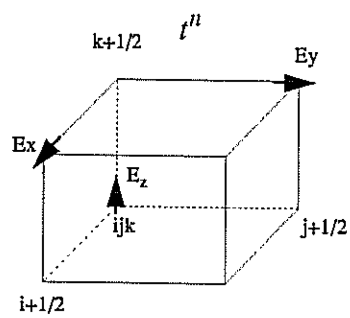
H components



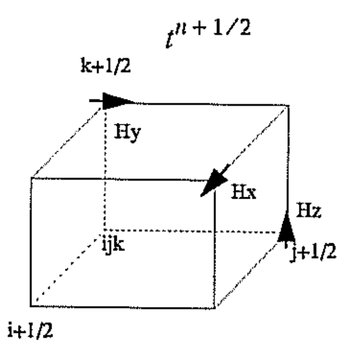
E components



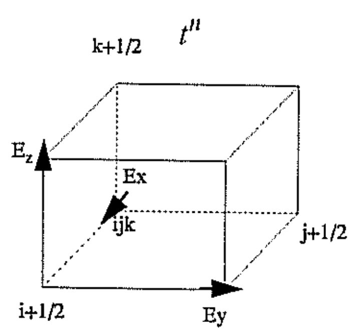
H components



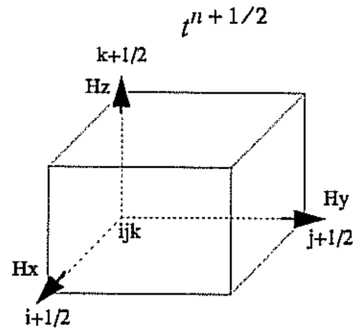
E components



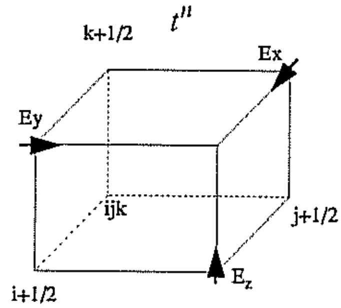
H components



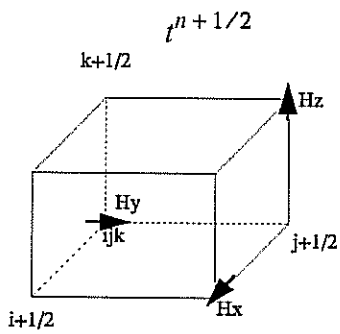
E components



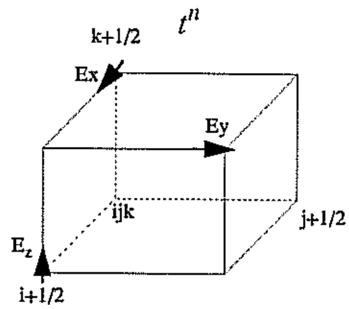
H components



E components



H components



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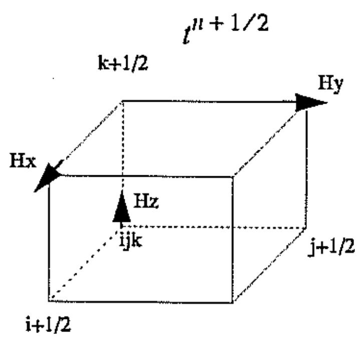
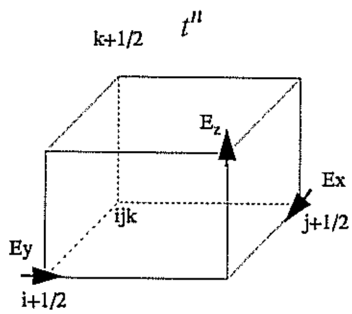
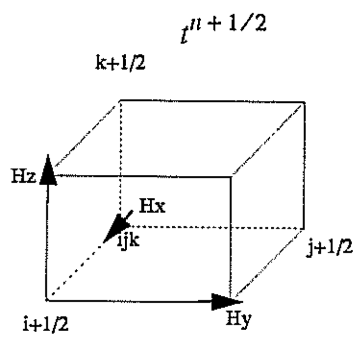


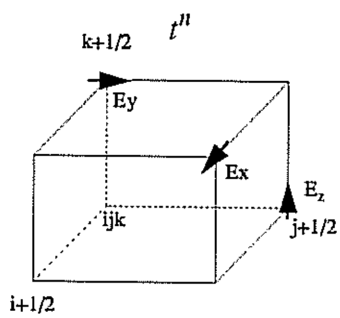
H components



E components



H components



E components